Introduction to Astronomy AA 201 Fall Semester 2019

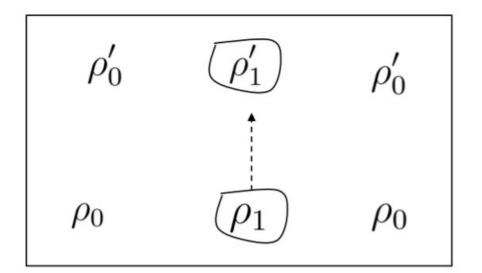
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Course webpage:

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Assume ho_1 is just slightly less than ho_0 Gas element rises by buoyancy -> new density is ho_1'

If $\rho_1'<\rho_0'$, it will keep rising, thus being unstable. Now, let's assume pressure equilibrium, so that $P_1'=P_0'$. Then $T_1'>T_0'$, so since the underdense pocket rises, heat is transported upwards, as long as it doesn't start transferring heat to its surroundings.

So condition for gas to keep rising is:

$$\frac{d\rho_1}{dr} < \frac{d\rho_0}{dr} \quad \text{(both } < 0\text{)}$$

If adiabatic (rises without transferring heat to surroundings), the adiabatic gas law for an ideal gas is obeyed:

$$P_1V_1^{\gamma} = Const$$

(what should γ be in a stellar interior?)

or

$$P_1 = k \rho_1^{\gamma}$$

Then
$$\frac{dP_1}{dr} = k\gamma \rho_1^{\gamma-1} \frac{d\rho_1}{dr} = k\gamma \frac{\rho_1^{\gamma}}{\rho_1} \frac{d\rho_1}{dr} = \gamma \frac{P_1}{\rho_1} \frac{d\rho_1}{dr}$$

So
$$\frac{d\rho_1}{dr} = \frac{1}{\gamma} \frac{\rho_1}{P_1} \frac{dP_1}{dr} < \frac{d\rho_0}{dr}$$
 (1)

We still have the ideal gas law: $P = \frac{\rho kT}{}$

$$P = \frac{\rho kT}{\mu m_H}$$

So we'll write:
$$\frac{dP_0}{dr} = \frac{P_0}{\rho_0} \frac{d\rho_0}{dr} + \frac{P_0}{T_0} \frac{dT_0}{dr}$$

Solve for $d
ho_0/dr$, substitute in (1). Assuming pressure equilibrium,

$$P_1 = P_0, \frac{dP_1}{dr} = \frac{dP_0}{dr}$$

and now assume that the initial densities were nearly equal, so approximately $\rho_1 = \rho_0$.

Then

$$\frac{dT_0}{dr} < \left(1 - \frac{1}{\gamma}\right) \frac{T_0}{P_0} \frac{dP_0}{dr}$$

But
$$\frac{dP_0}{dr} = -\rho_0 g = -\rho_0 \frac{GM_r}{r^2}$$
 and $\frac{T_0}{P_0} = \frac{\mu m_H}{\rho_0 k}$

$$\frac{dT_0}{dr} < -\left(1 - \frac{1}{\gamma}\right) \frac{\mu m_H}{k} \frac{GM_r}{r^2} = \frac{dT}{dr}|_{ad}$$

condition for convection

(Remember, dT/dr < 0)

Since the temperature gradients are negative, then if $|\frac{dT}{dr}| > |\frac{dT}{dr}|_{ad}$ bubbles will rise, and heat will be carried by convection.

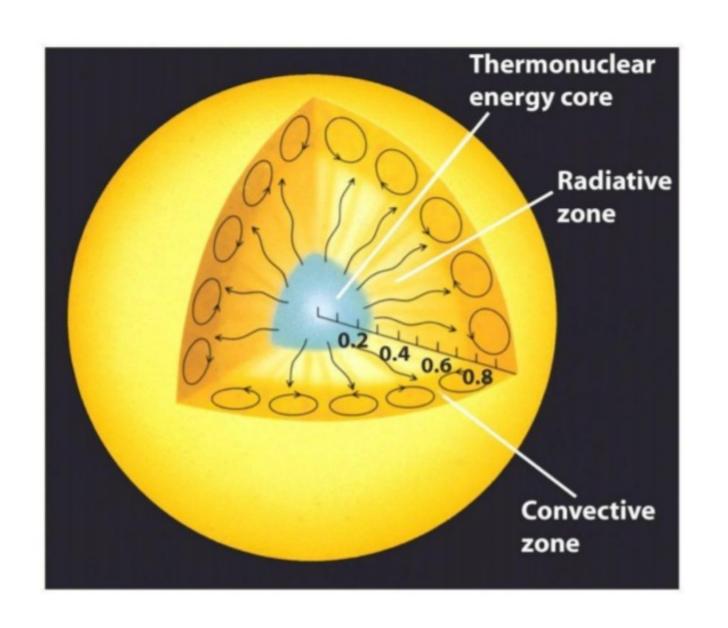
Stops when bubble can finally cool (i.e. is no longer adiabatic), e.g. near the top of the Sun's atmosphere.

With the convection condition in this form, we can see when convection is likely, namely when

1.
$$\overline{\mathcal{K}}$$
 is large (so $|\frac{dT}{dr}|$ is large)

2.
$$\frac{L_r}{4\pi r^2}$$
 is high, deep in cores of massive stars (so $|\frac{dT}{dr}|$ is large)

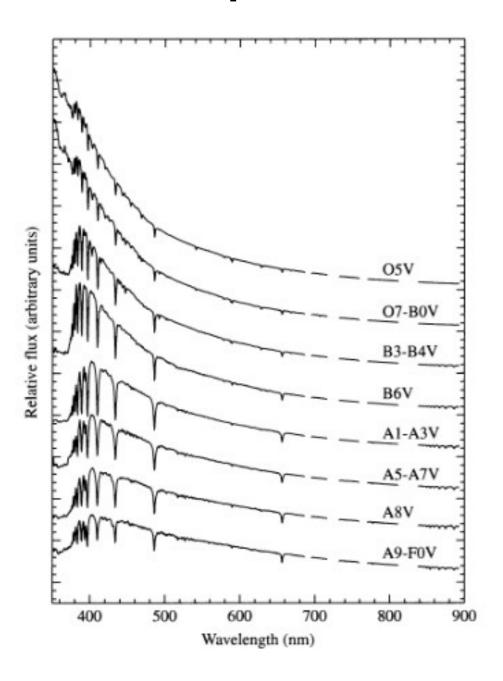
3.
$$g=\frac{GM_r}{r^2}$$
 is low (so $|\frac{dT}{dr}|_{ad}$ is low)

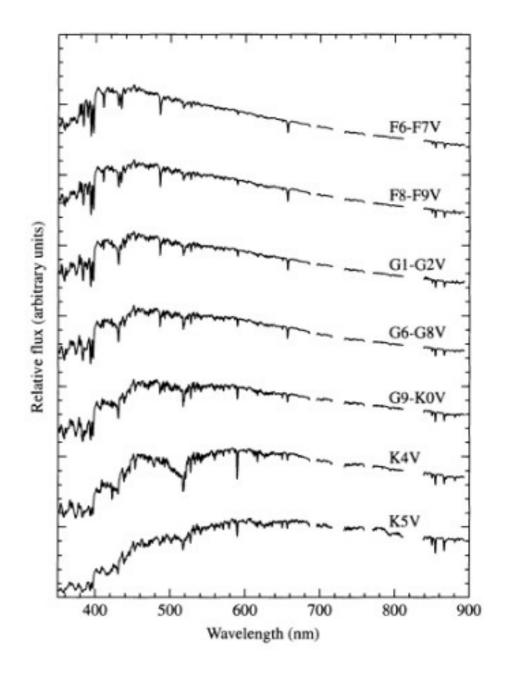


Spectral Type	Characteristics
0	Hottest blue-white stars with few lines Strong He II absorption (sometimes emission) lines. He I absorption lines becoming stronger.
В	Hot blue-white He I absorption lines strongest at B2. H I (Balmer) absorption lines becoming stronger.
Α	White Balmer absorption lines strongest at A0, becoming weaker later. Ca II absorption lines becoming stronger.
F	Yellow-white Ca II lines continue to strengthen as Balmer lines continue to weaken. Neutral metal absorption lines (Fe I, Cr I).
G	Yellow Solar-type spectra. Ca II lines continue becoming stronger. Fe I, other neutral metal lines becoming stronger.
K	Cool orange Ca II H and K lines strongest at K0, becoming weaker later. Spectra dominated by metal absorption lines.
М	Cool red Spectra dominated by molecular absorption bands, especially titanium oxide (TiO) and vanadium oxide (VO). Neutral metal absorption lines remain strong.
L	Very cool, dark red Stronger in infrared than visible. Strong molecular absorption bands of metal hydrides (CrH, FeH), water (H ₂ O), carbon monoxide (CO), and alkali metals (Na, K, Rb, Cs). TiO and VO are weakening.
Т	Coolest, Infrared Strong methane (CH ₄) bands but weakening CO bands.

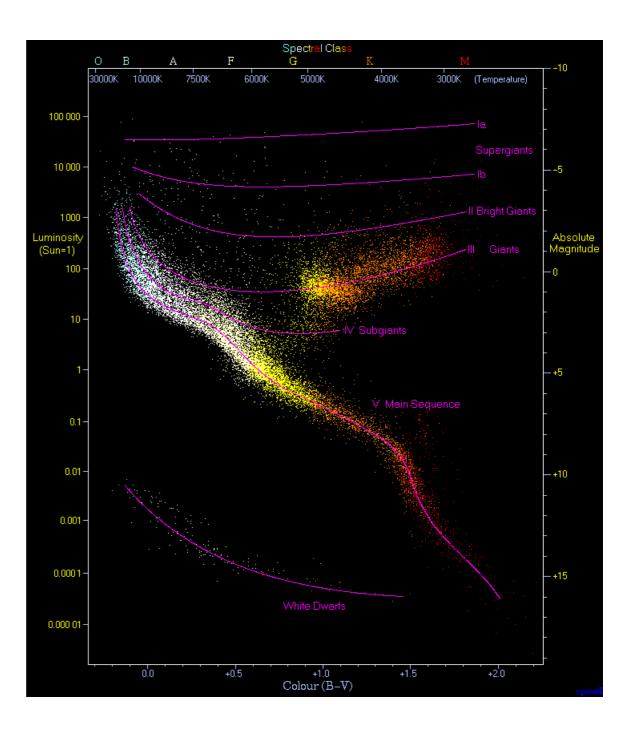
Stellar Spectral Classification Scheme

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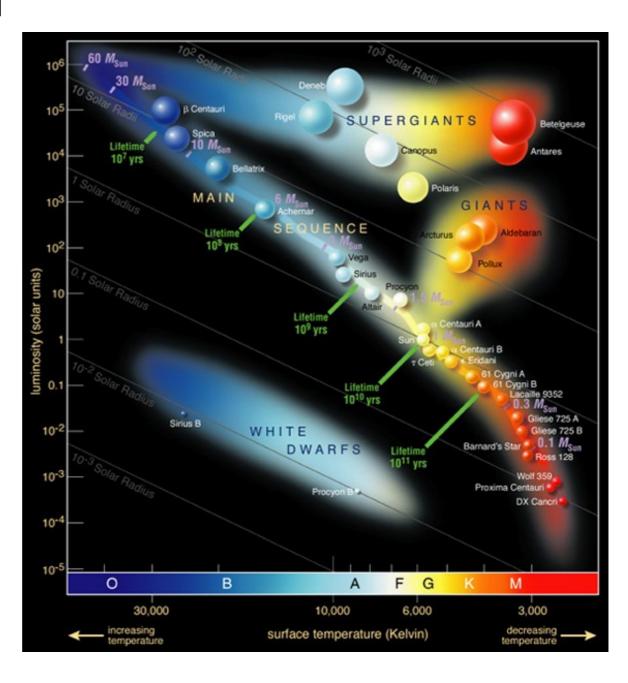


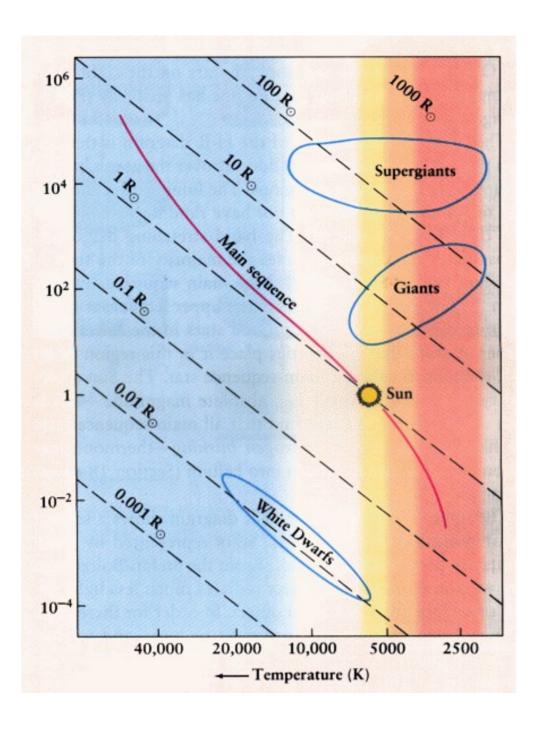


H-R diagram

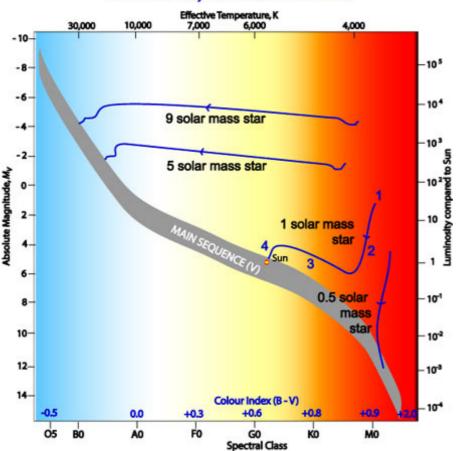


H-R diagram





Theoretical Hayashi Tracks of Protostars



Evolutionary Tracks off the Main Sequence

