

Galactic and Extragalactic Astronomy

AA 472/672

Spring Semester

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- **Oldest galaxies in universe**

- most of their stars formed early in universe;
- the galaxy may have grown or changed since early universe

- **Appear simple but are complex**

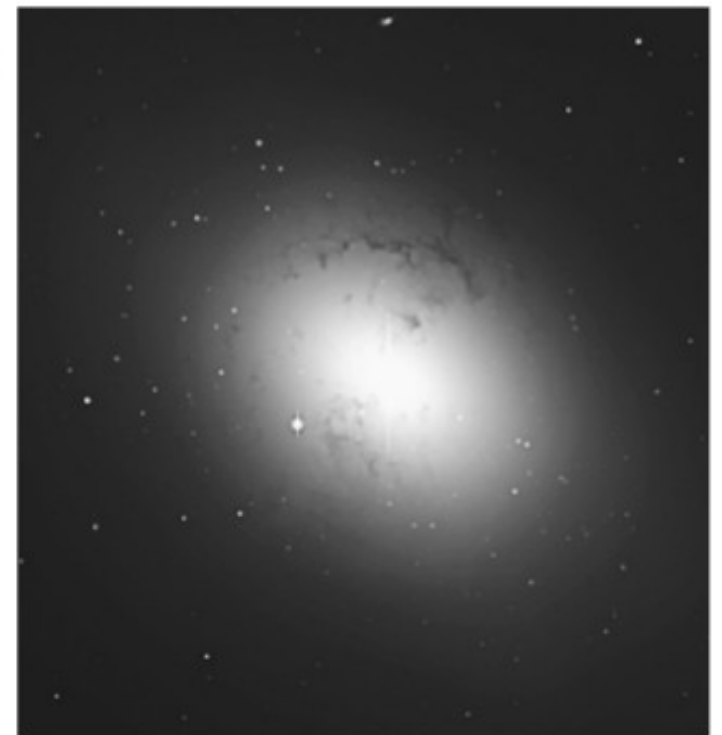


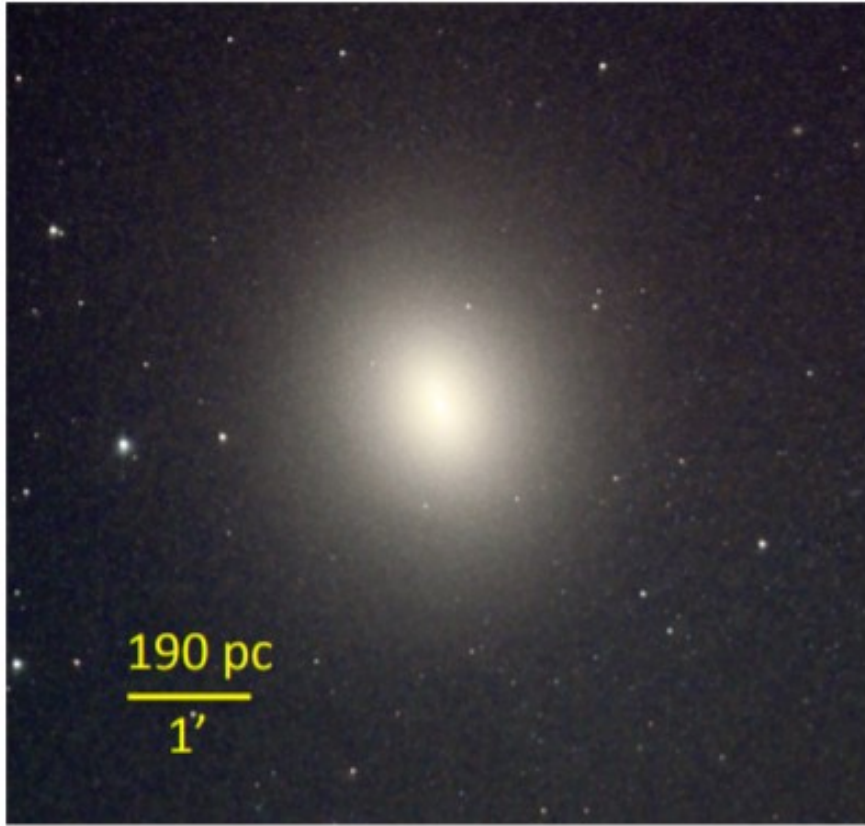
A typical elliptical galaxy

Characteristics of elliptical galaxies

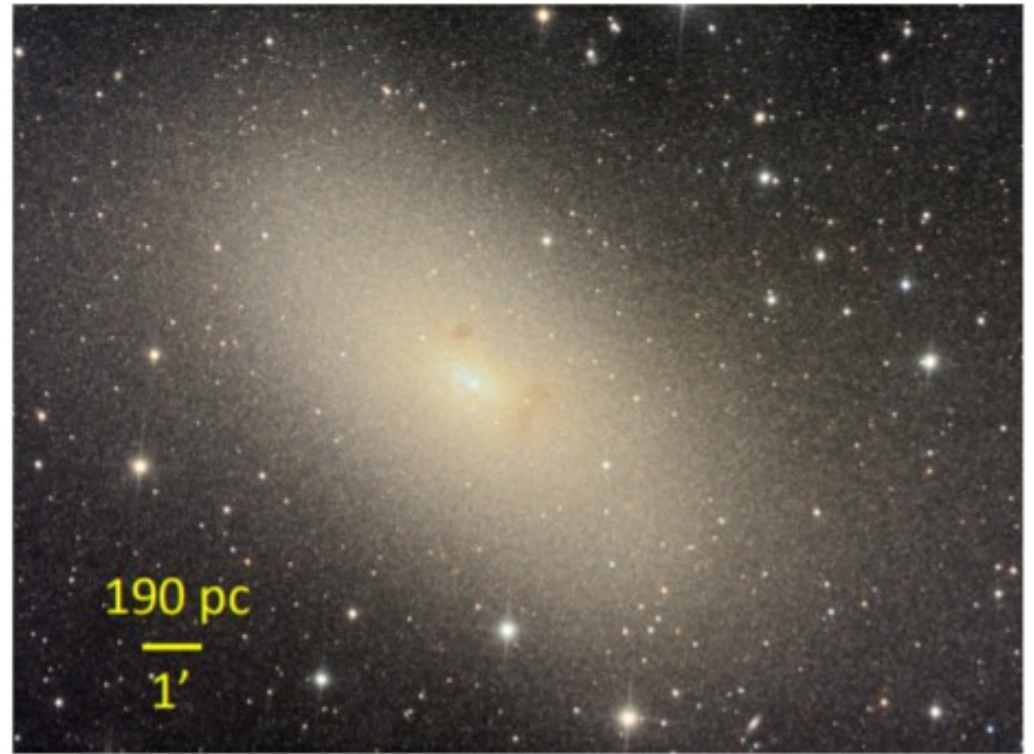
- Little or no star formation
- Little or no dust or cold gas
- Little or no substructure within galaxy
- Isophote shapes nearly ellip/cal

Real story is much more complex

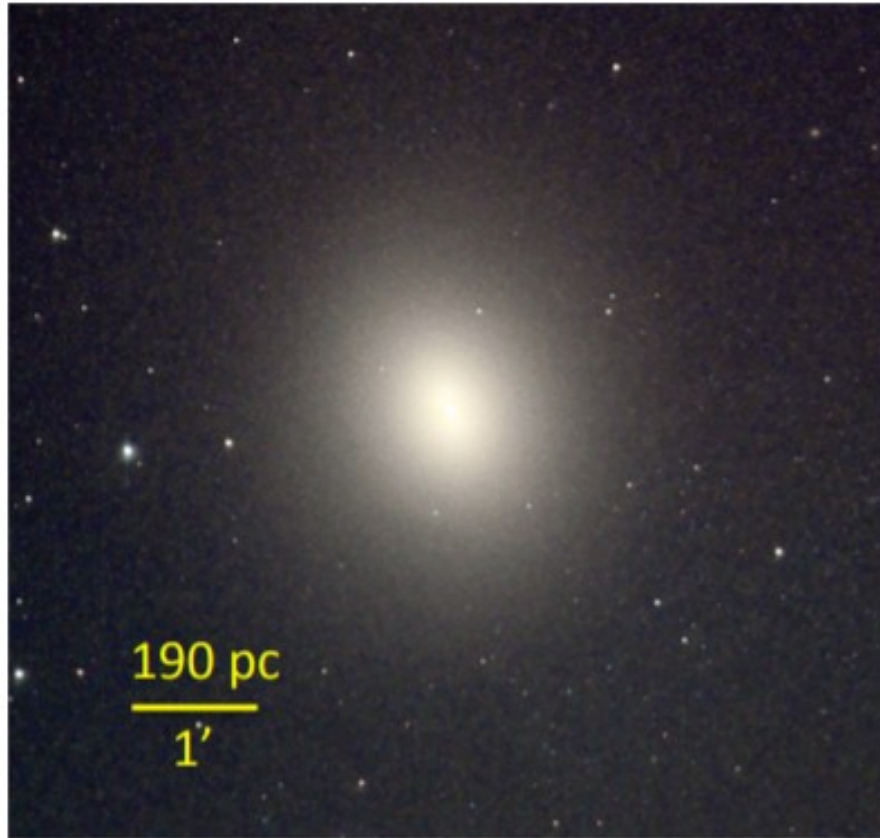




Elliptical M32



Dwarf "elliptical" NGC 205



Elliptical M32

Compact, high central stellar density of stars

Little or no gas & star formation



Dwarf “elliptical” NGC 205

Not compact, low central surface density of stars

Little or no gas & star formation

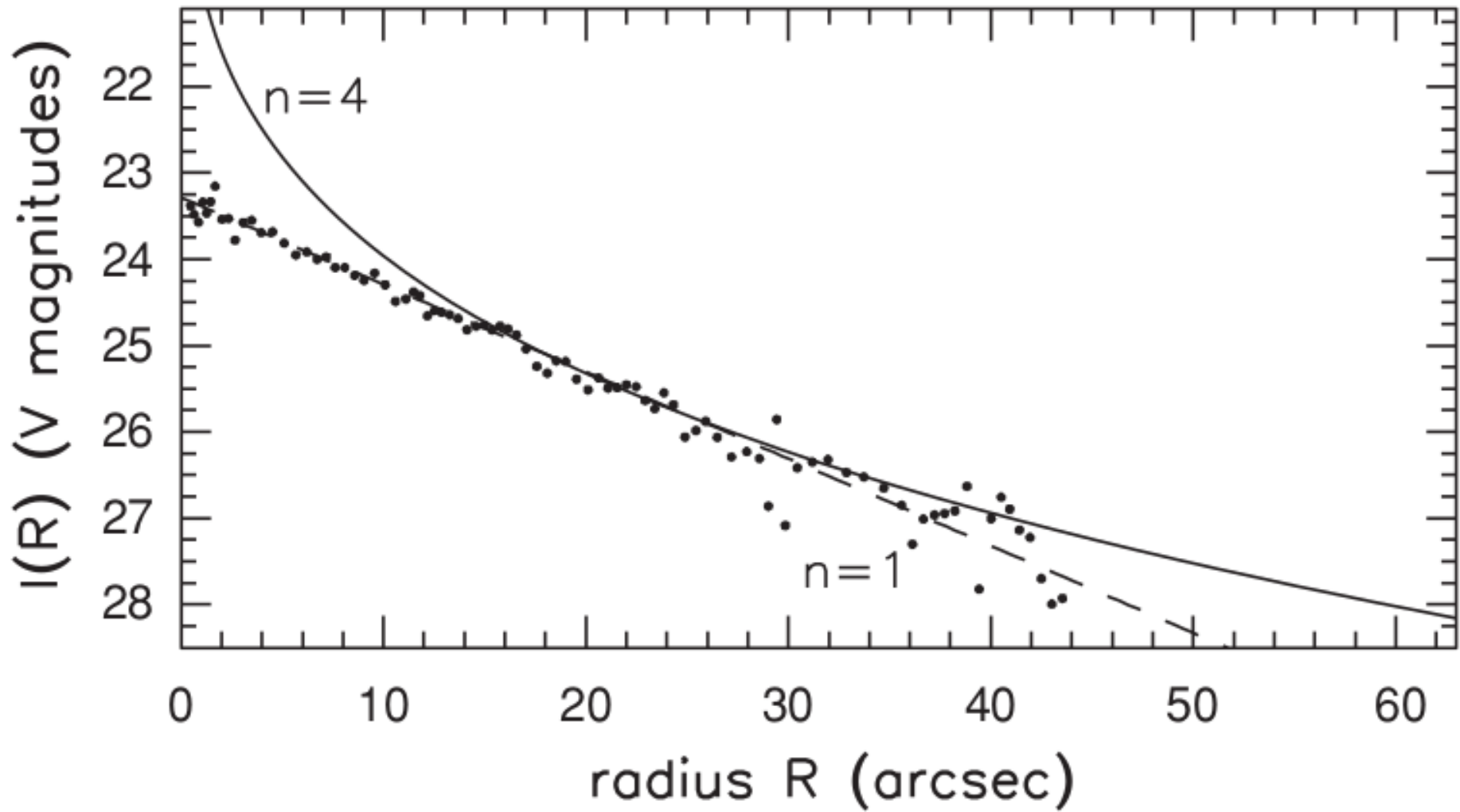
Brightness profile

$$I(R) = I_e \exp \left(-7.669 \left[(R/R_e)^{1/4} - 1 \right] \right)$$

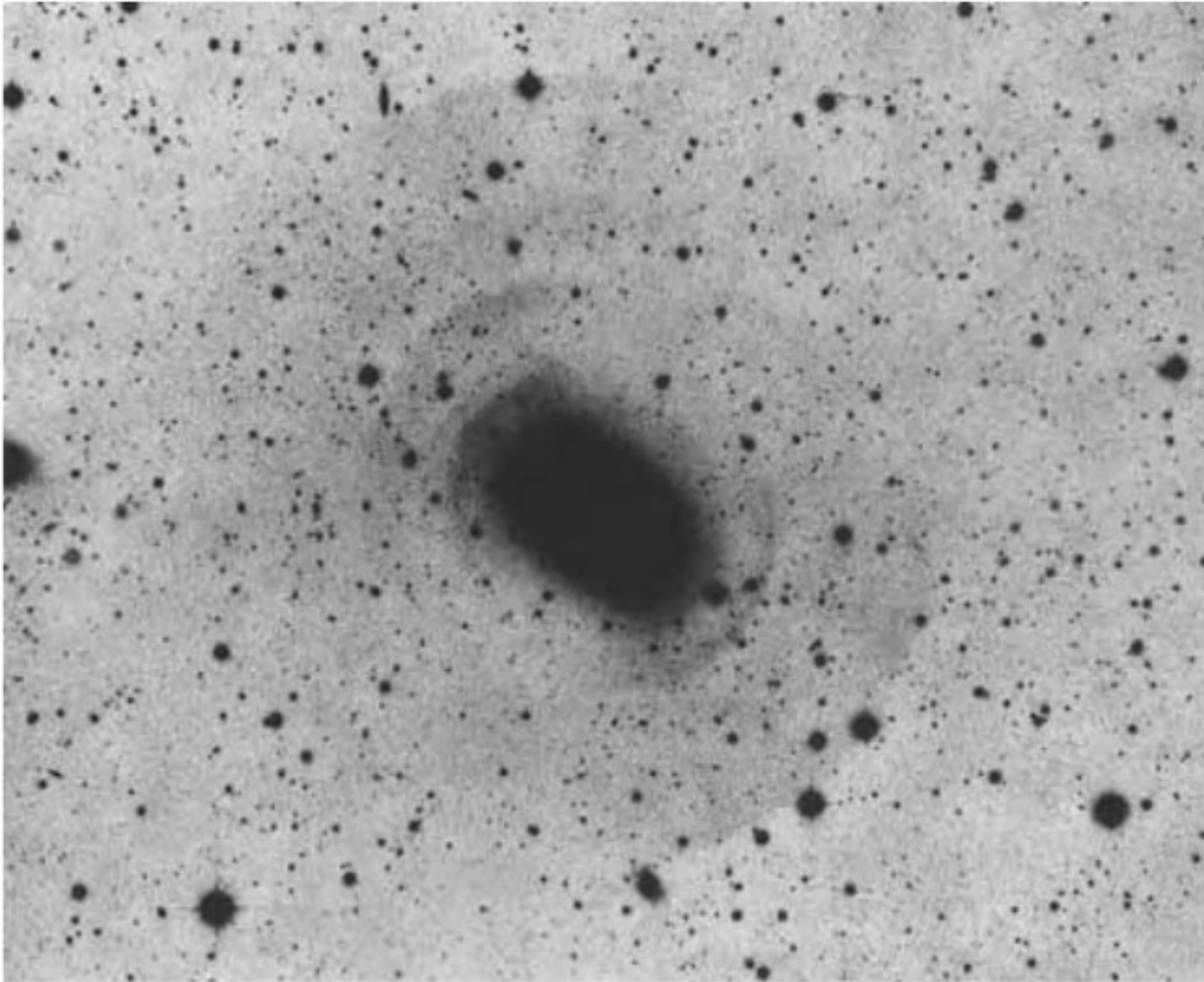
$I(R)$ → surface brightness

R_e → effective radius defined such that half of the luminosity is emitted from within R_e

de Vaucouleurs Profile



galaxy VCC753 in the Virgo cluster



NGC 3923

Arclike structures → signatures of merger or tidal stripping

Exponential disk:

$$I(r) = I(0) \exp(-r/r_d)$$

deVaucouleurs $r^{1/4}$ bulge law: $I(r) = I(r_{\text{eff}}) \exp \{-7.67[(r/r_{\text{eff}})^{1/4}-1]\}$

Sersic law:

$$I(r) = I(r_{\text{eff}}) \exp \{-b_n[(r/r_{\text{eff}})^{1/n}-1]\}$$

n = Sersic index

b_n chosen to make r_{eff} the effective radius (encloses $\frac{1}{2}$ the light)

$$b_n = 1.999n - 0.327 \text{ for } n > 1$$

$n = 1-4$ typically

If $n=1$ exponential (all disk) disks of spirals, S0s, dwarf Es

If $n=4$ deVaucouleurs $r/4$ law (all bulge) giant E's, globular clusters

$1 < n < 4$ bulges of spirals and S0s (higher n for large L bulges)

If $n < 2$ for entire spiral or S0: small bulge-disk ratio

If $n > 2$ for entire spiral or S0: large bulge-disk ratio

Advantage of Sersic law: *can describe entire profile shape with just 1 number n*

Problem 6.1 Show that the $R^{1/4}$ formula yields a total luminosity

$$L = \int_0^\infty 2\pi R I(R) dR = 8! \frac{e^{7.67}}{(7.67)^8} \pi R_e^2 I(R_e) \approx 7.22 \pi R_e^2 I(R_e). \quad (6.2)$$

(Remember that $\int_0^\infty e^{-t} t^7 dt = \Gamma(8) = 7!$) Use a table of incomplete Γ functions to show that half of this light comes from within radius R_e .

Properties characterizing E's

- Little or no star formation
- Little or no dust or cold gas
- Little or no substructure within galaxy
- Isophote shapes nearly elliptical

If you use just these properties, you include both “real ellipticals” as well as dwarf galaxies that are not true ellipticals

Properties characterizing E's

- Little or no star formation
- Little or no dust or cold gas
- Little or no substructure within galaxy
- Isophote shapes nearly elliptical
- Radial light distribution: $n \cong 4$

The Faber–Jackson relation and the fundamental plane

V_{rot} → rotational velocity

σ_v → velocity dispersion

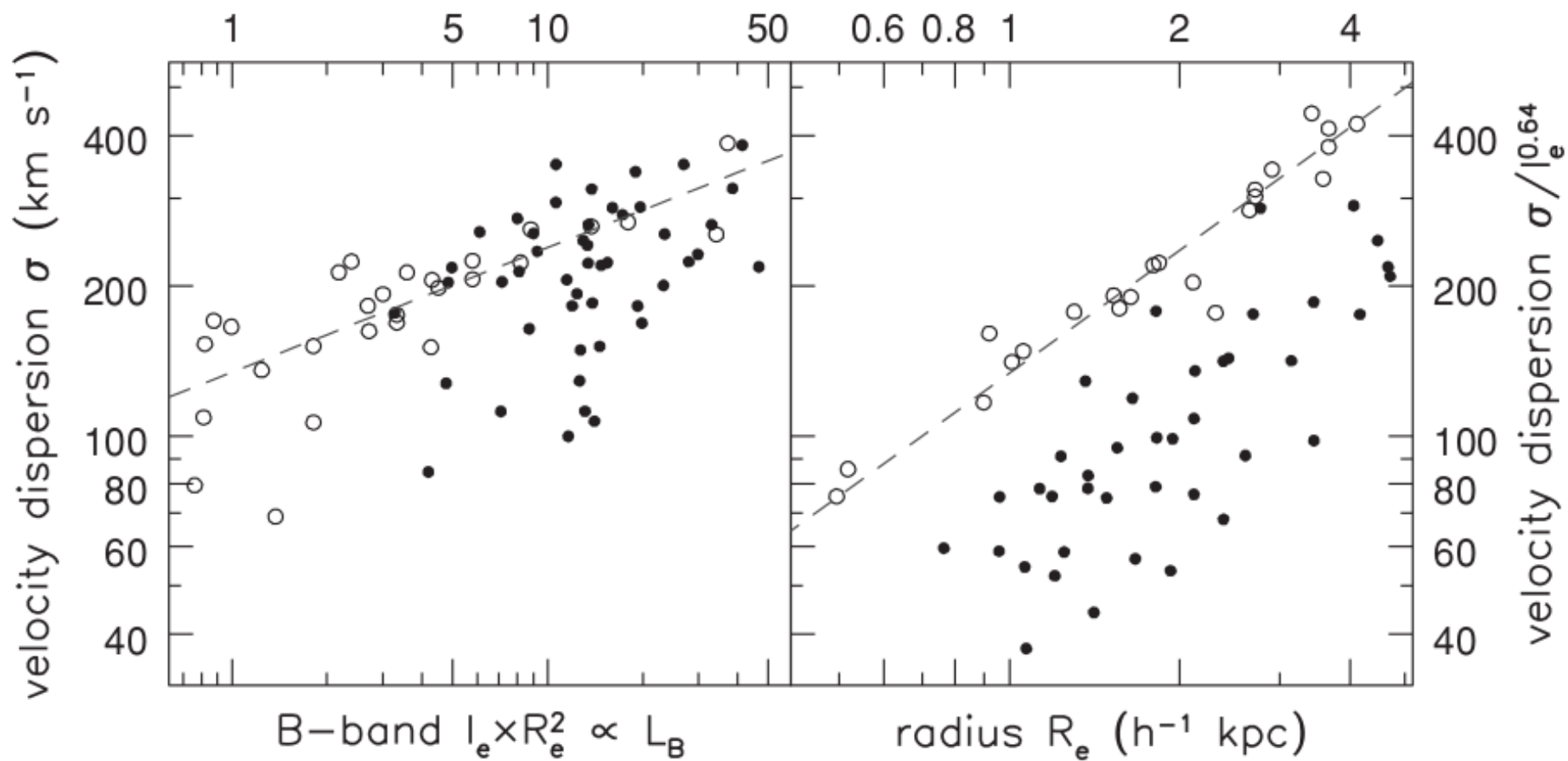
$$\frac{L_V}{2 \times 10^{10} L_{\odot}} \approx \left(\frac{\sigma}{200 \text{ km s}^{-1}} \right)^4.$$

Stars move faster in more luminous galaxies

At the centers of bright ellipticals, the dispersion can reach 500 km s^{-1} , while $\sigma \sim 50 \text{ km s}^{-1}$ in the least luminous objects

Velocity measurement using 21 cm line

Line widths give estimates of velocity dispersion



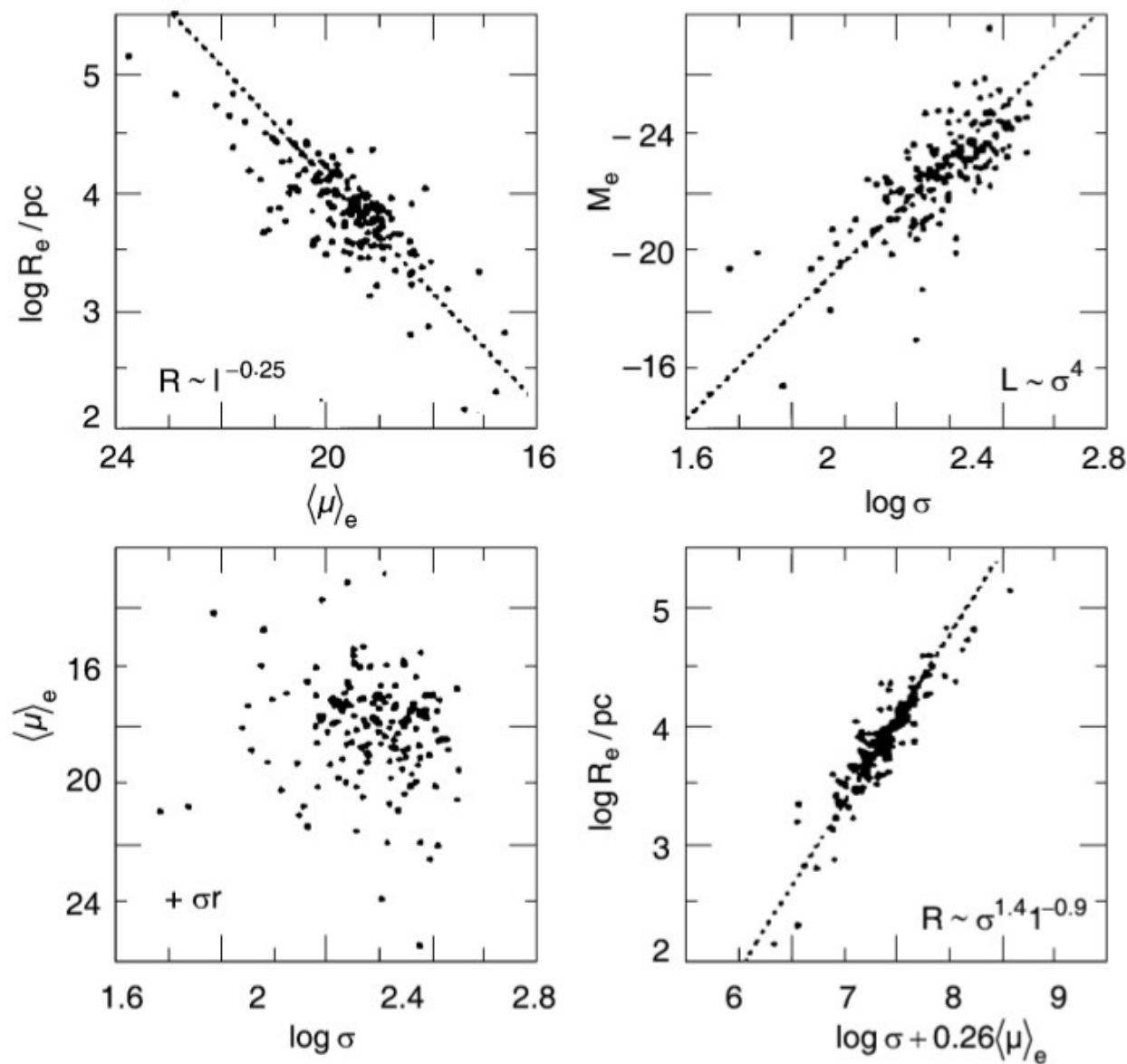
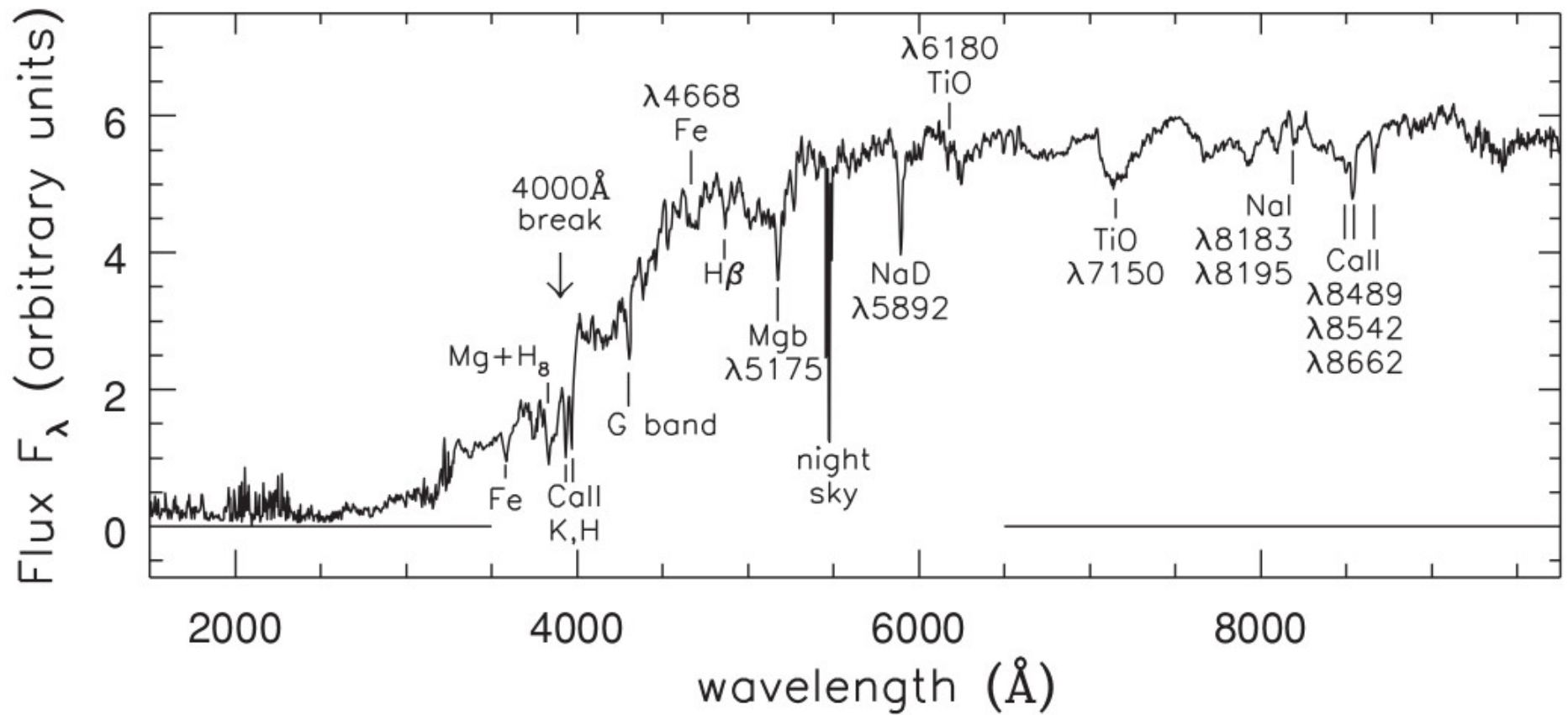


Fig. 3.23. Projections of the fundamental plane onto different two-parameter planes. Upper left: the relation between radius and mean surface brightness within the effective radius. Upper right: Faber–Jackson relation. Lower left: the relation between mean surface brightness and velocity dispersion shows the fundamental plane viewed from above. Lower right: the fundamental plane viewed from the side – the linear relation between radius and a combination of surface brightness and velocity dispersion



Spectrum of a typical elliptical galaxy

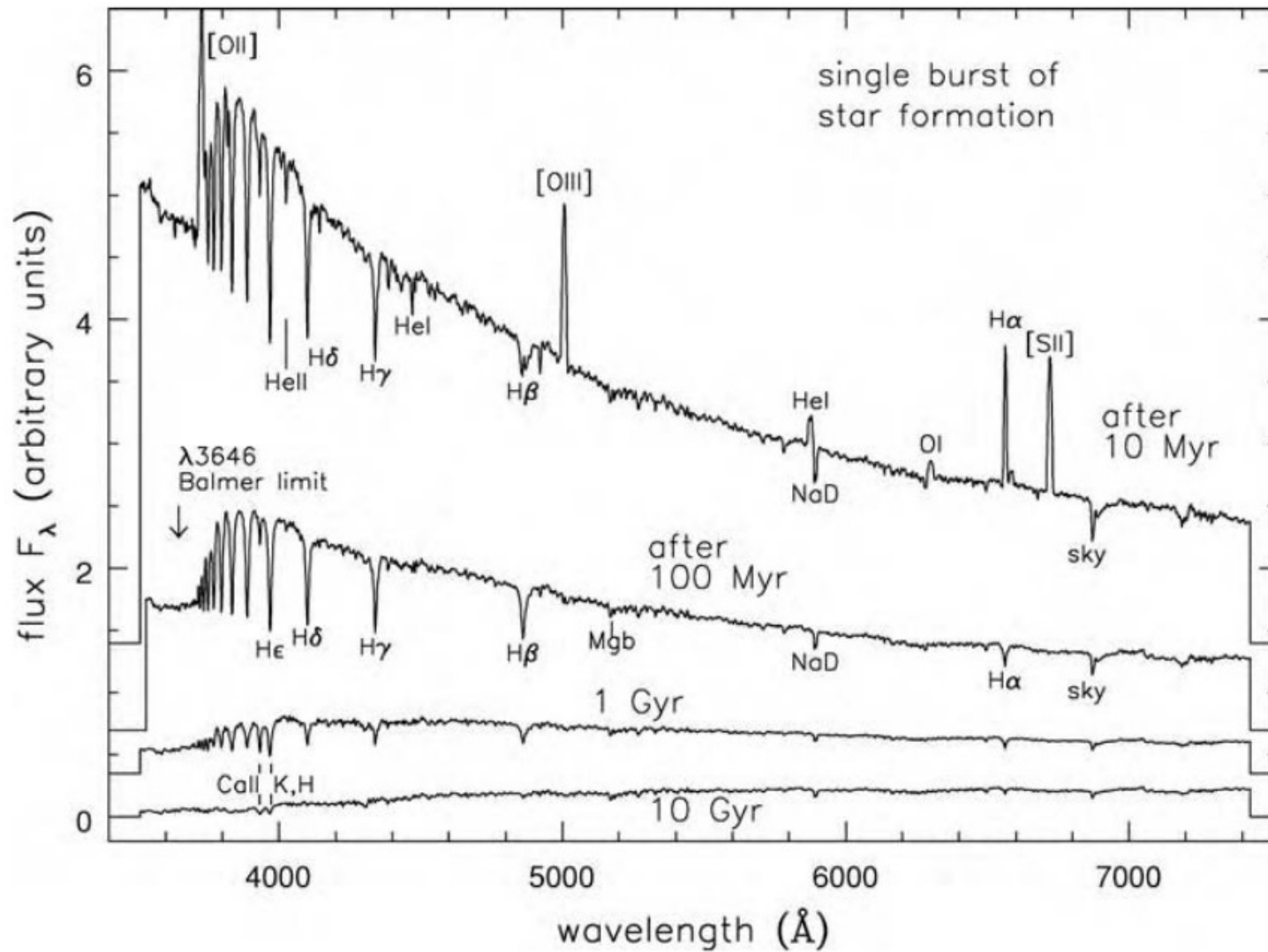
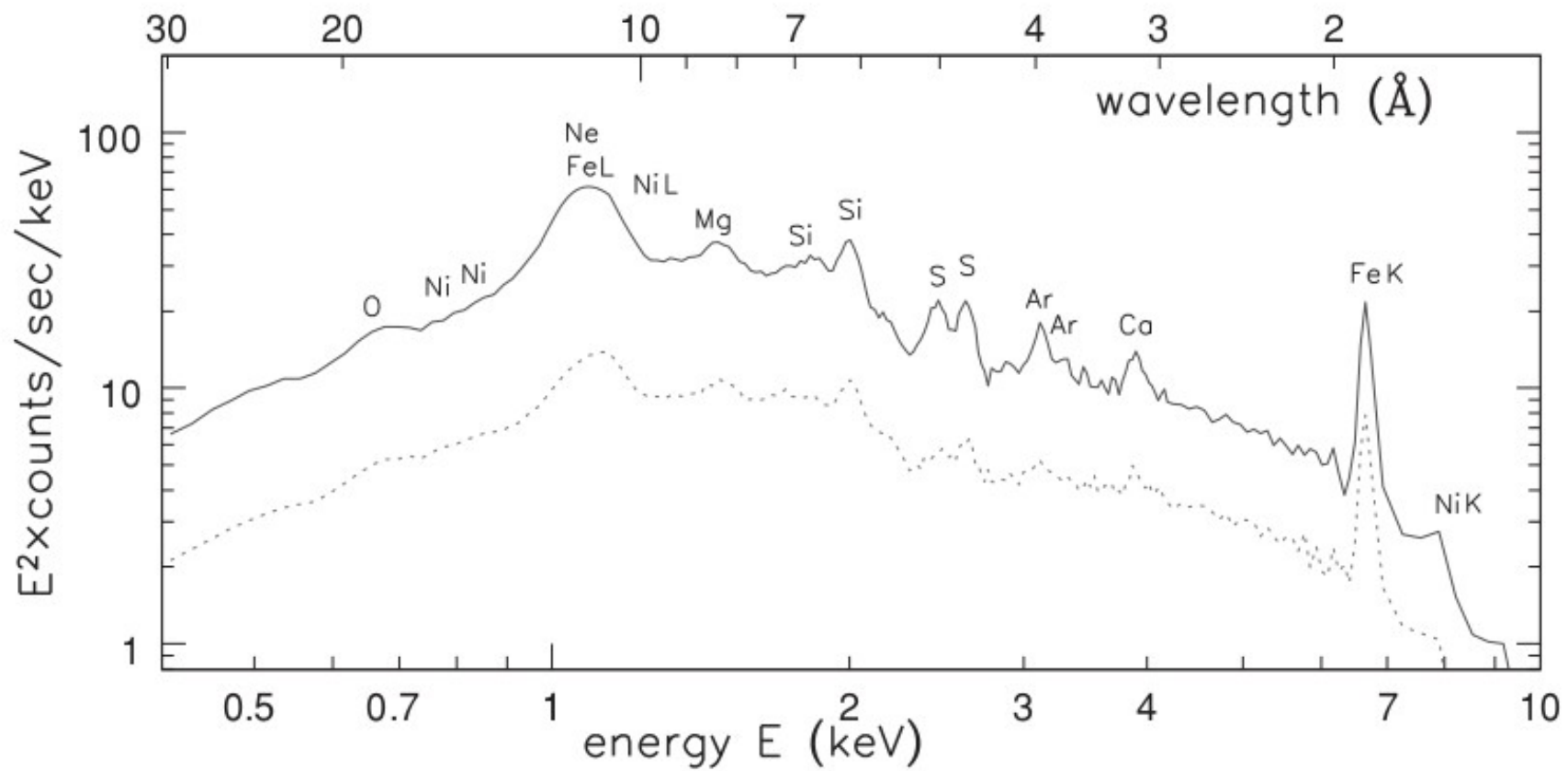


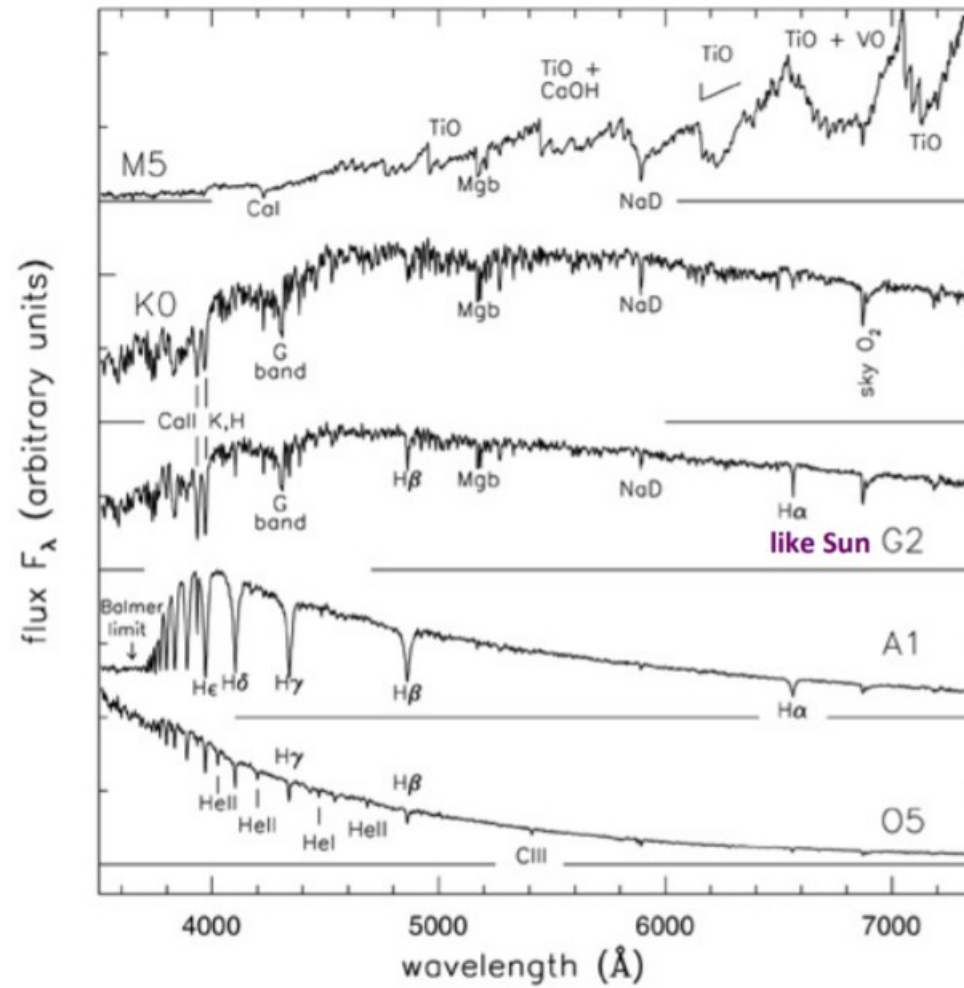
Fig. 6.18. Spectra for a ‘galaxy’ that makes its stars in a 10^8 yr burst, all plotted to the same vertical scale. Emission lines of ionized gas are strong 10 Myr after the burst ends; after 100 Myr, the galaxy has faded and reddened, and deep hydrogen lines of A stars are prominent. Beyond 1 Gyr, the light dims and becomes slightly redder, but changes are much slower – B. Poggianti.



H

Hot gas
M 87

Spectra of main-sequence stars



Compare the spectra of K giant star & S0 galaxy

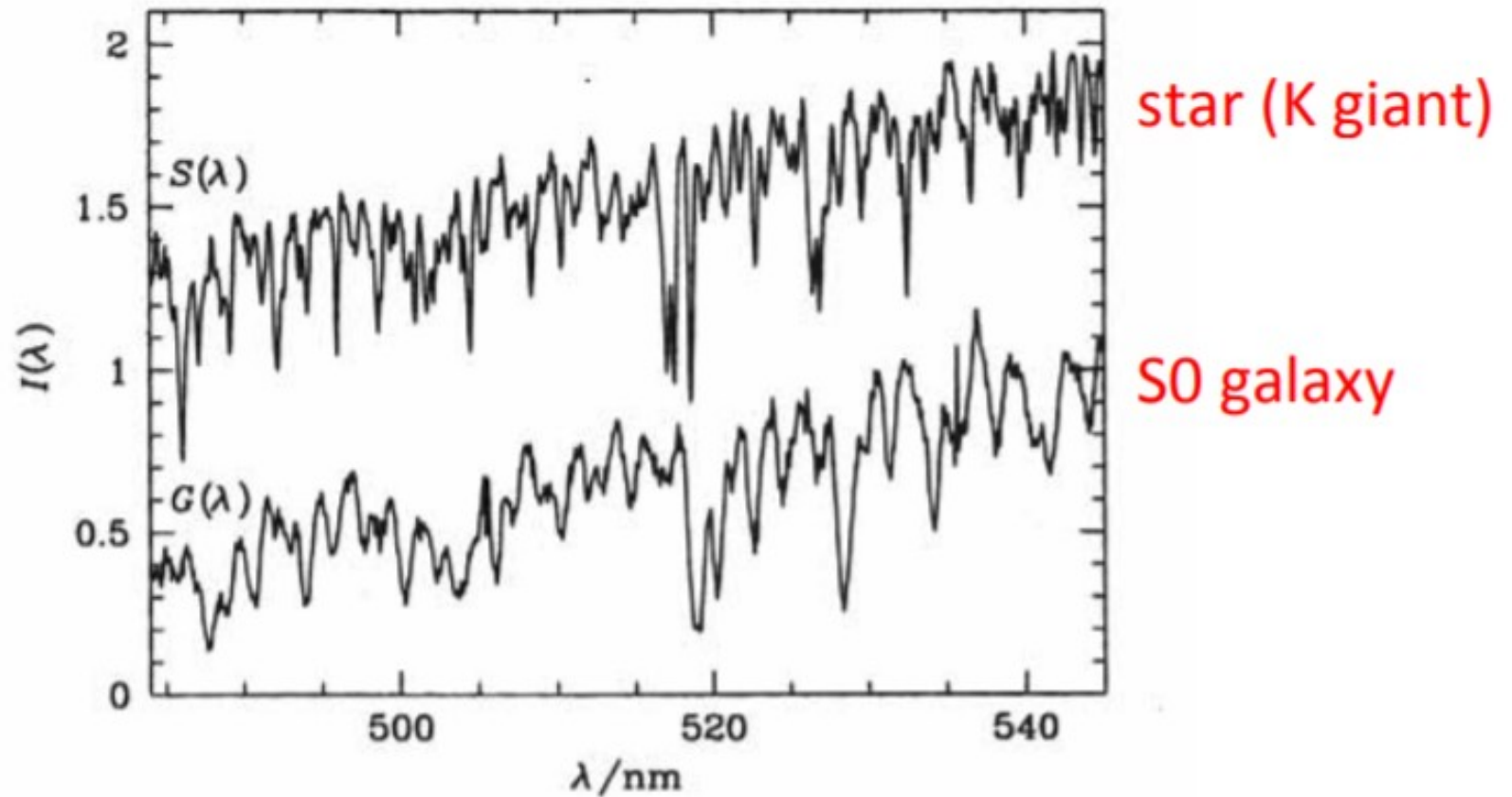


Figure 11.1 Spectra of a K0 giant star (S) and the center of the lenticular galaxy NGC 2549 (G). These data cover a small part of the optical spectrum around the strong Mg b absorption feature at 518 nm.

Q: How are these spectra different?

Compare the spectra of K giant star & S0 galaxy

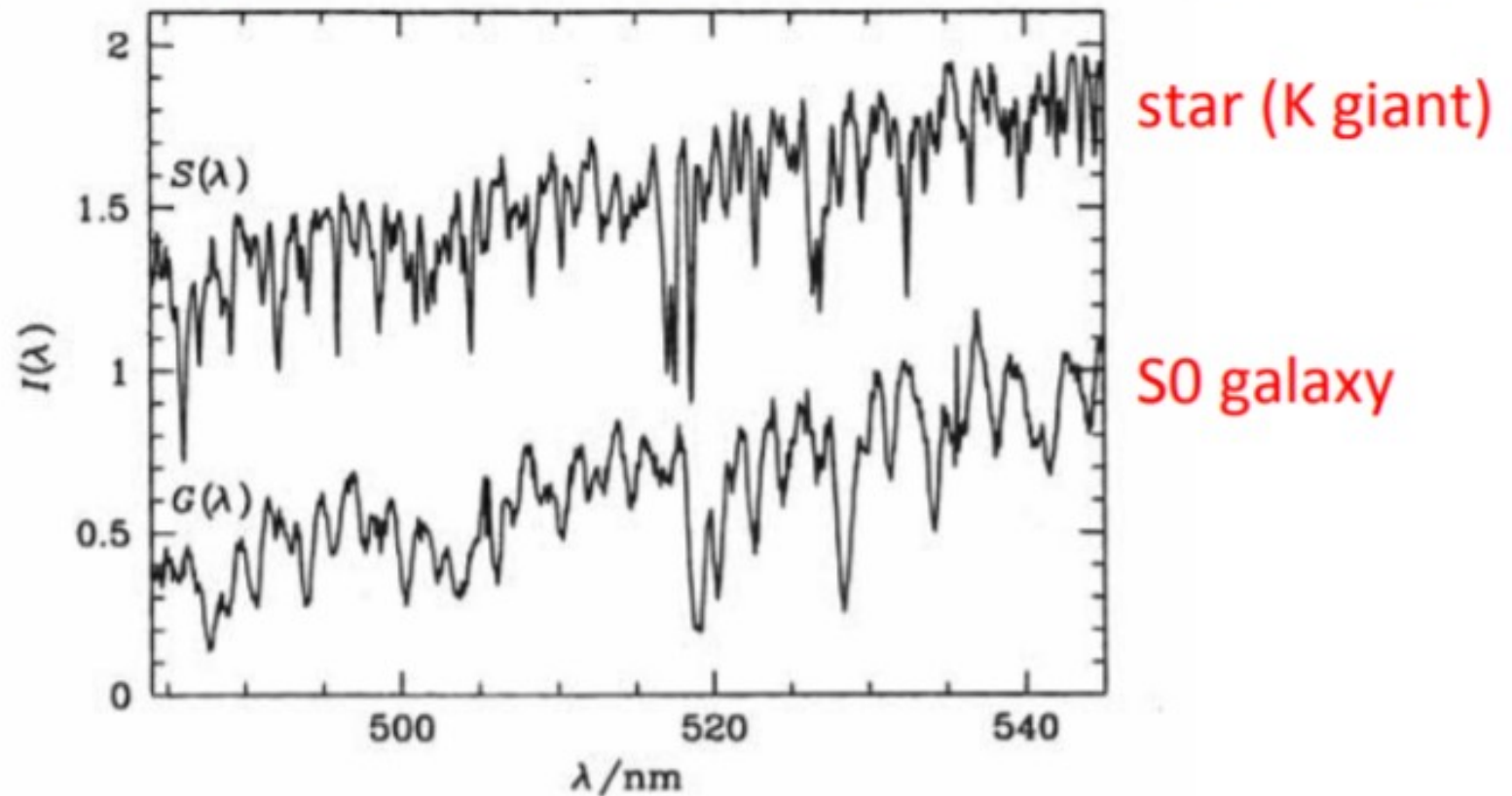


Figure 11.1 Spectra of a K0 giant star (S) and the center of the lenticular galaxy NGC 2549 (G). These data cover a small part of the optical spectrum around the strong Mg b absorption feature at 518 nm.

2 differences:

1. galaxy spectrum is redshifted wrt MW star (expansion of universe)

2. lines broader in galaxy due to velocity smearing

Introduction to kinematics for Ellipticals

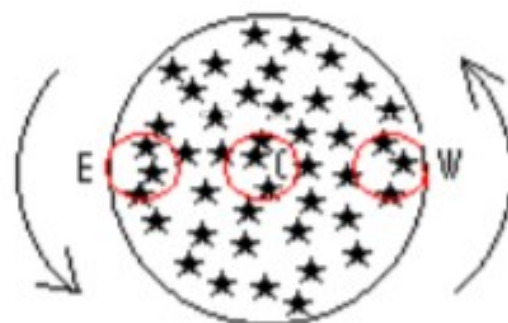
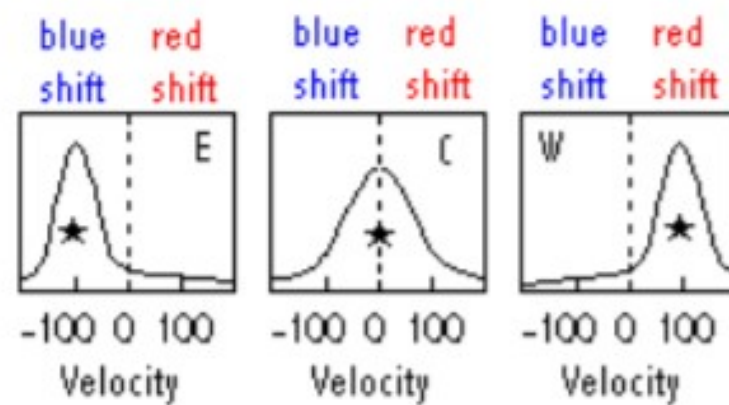
Ordered motions:

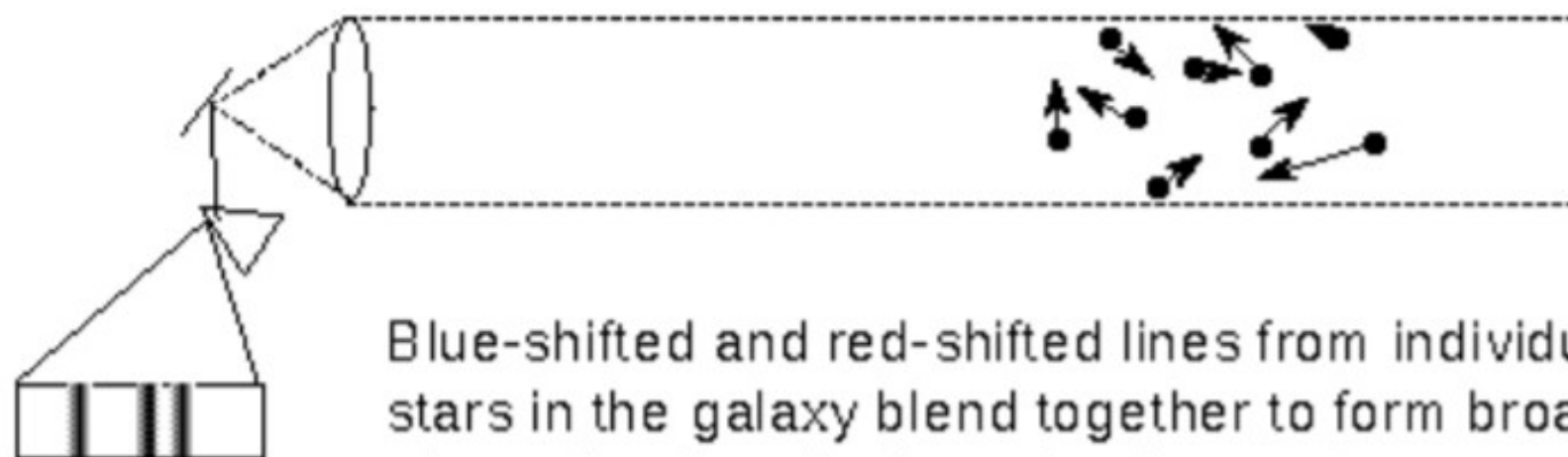
v : mean velocity $v = v_{\text{rot}} + v_{\text{noncirc}}$ *measured by peak or mean of line*

Disordered motions:

σ : velocity dispersion, *measured by linewidth*

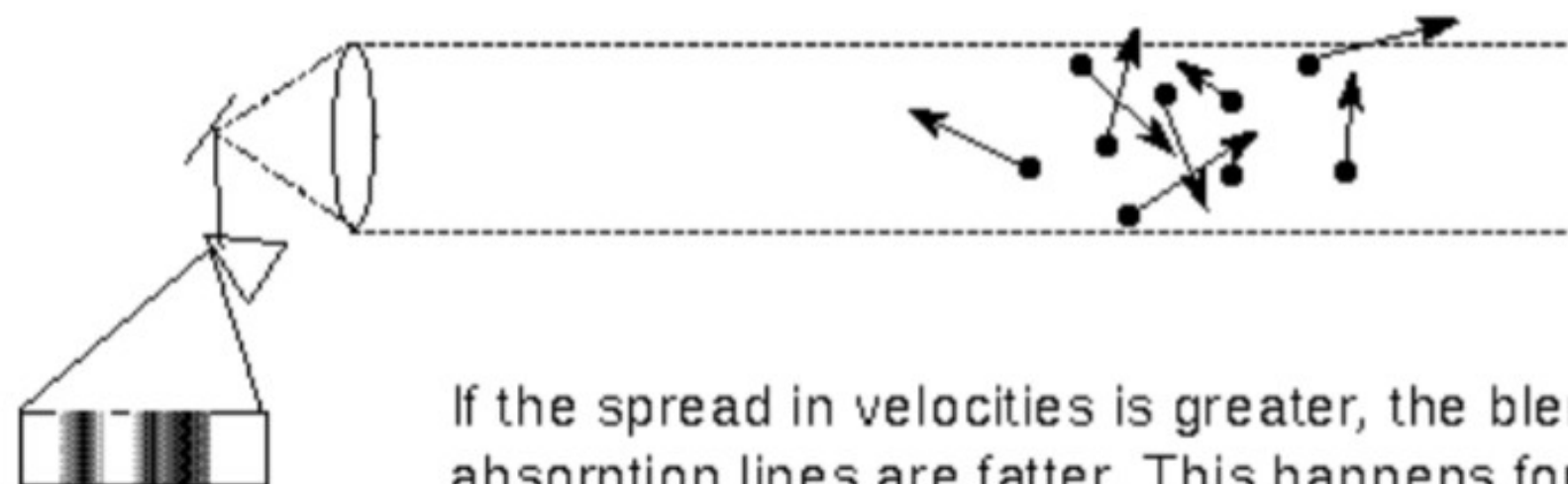
The ratio v/σ is used to compare the relative importance of ordered and random motions





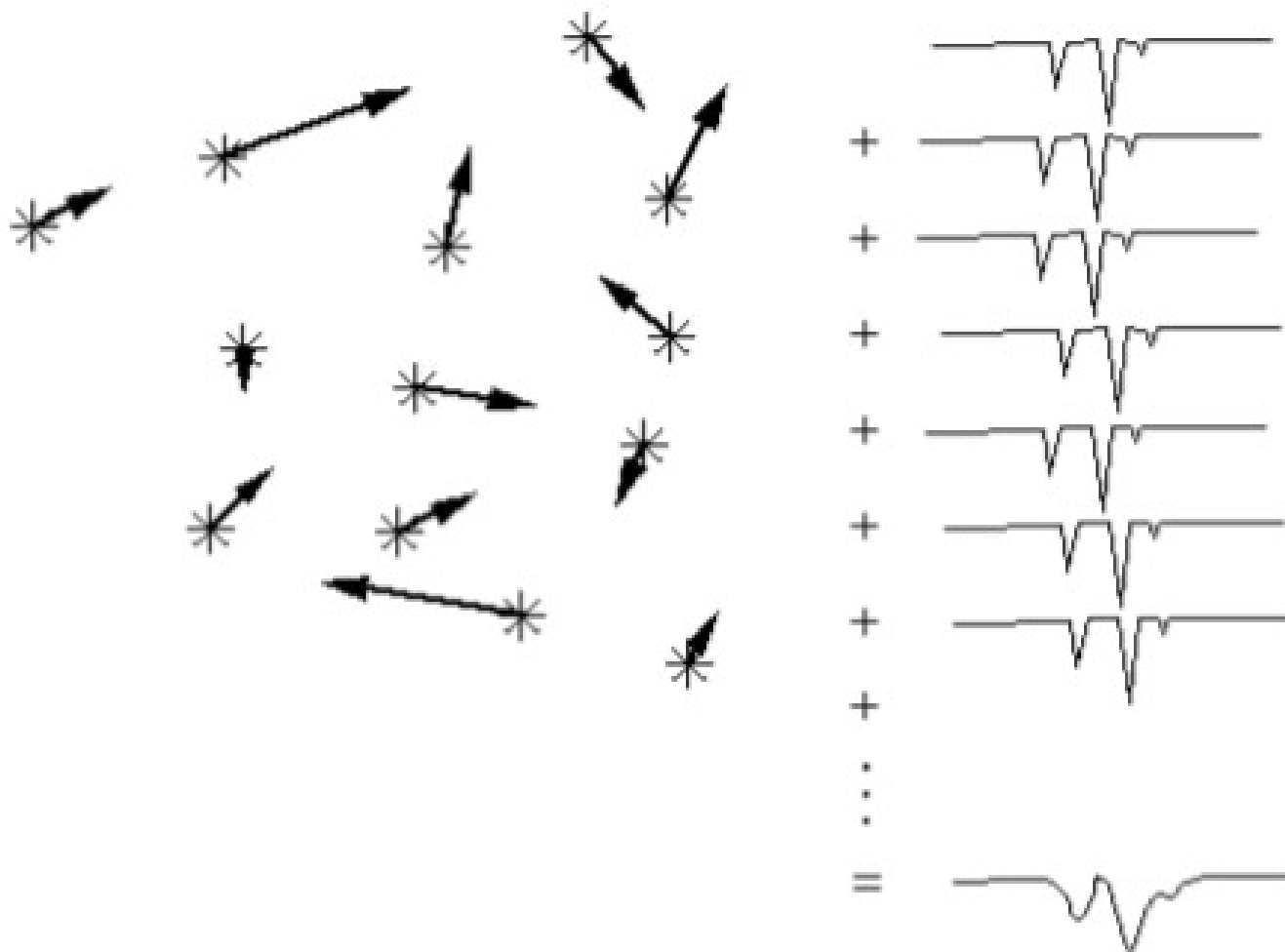
small σ

Blue-shifted and red-shifted lines from individual stars in the galaxy blend together to form broadened absorption lines in the galaxy's spectrum.



large σ

If the spread in velocities is greater, the blended absorption lines are fatter. This happens for the more massive and luminous galaxies.



LOSVD

Number
of stars



Velocity dispersion σ --
fit LOSVD with gaussian
(even if distribution is not gaussian!)

LOS velocities of stars

$$R_e \propto \sigma^{1.2} I_e^{-0.8}. \quad (6.19)$$

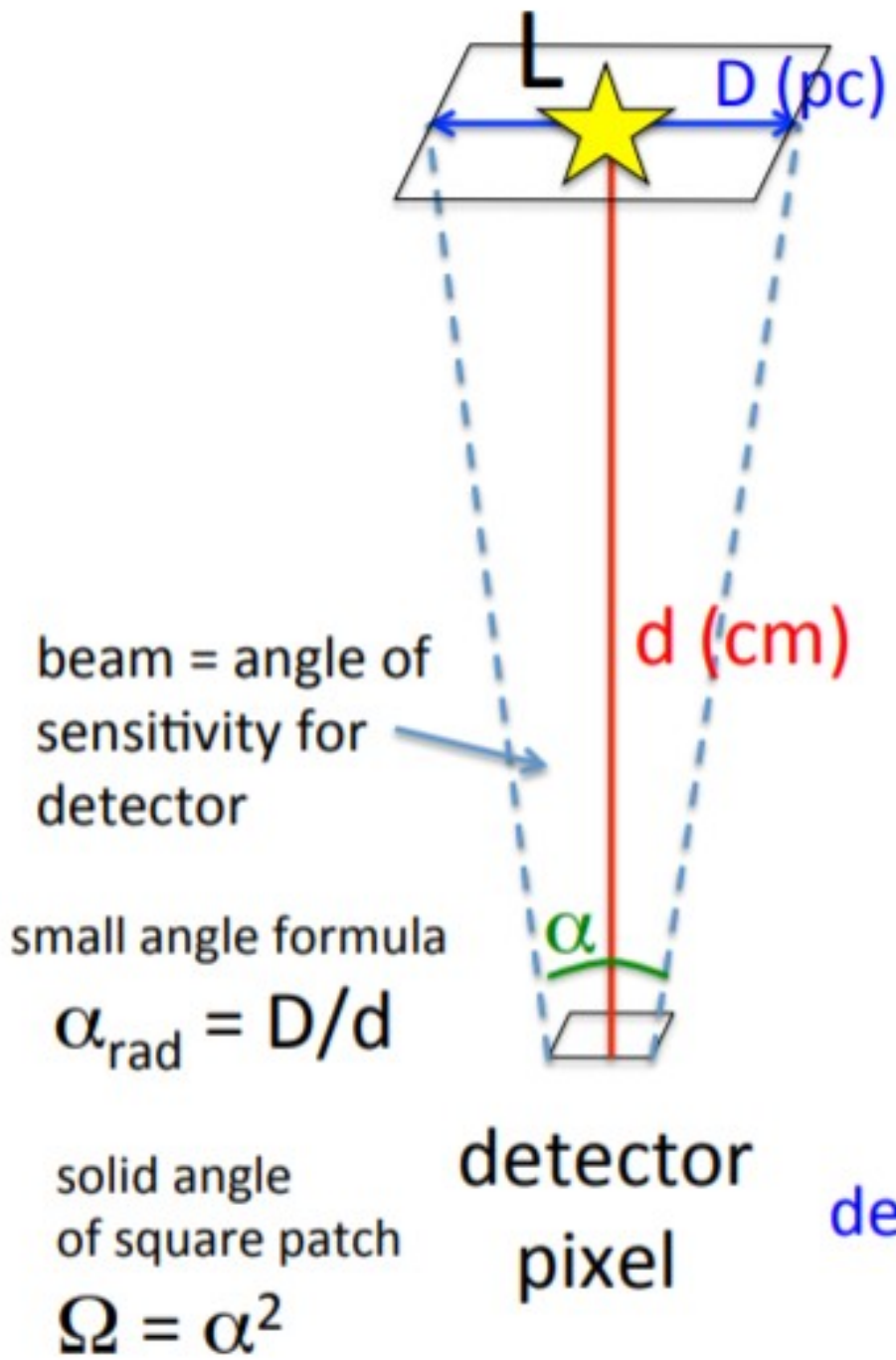
Problem 6.7 Assuming that the velocity dispersion σ and the ratio \mathcal{M}/L are roughly constant throughout the galaxy, and that no dark matter is present, show that the kinetic energy $\mathcal{KE} = 3\mathcal{M}\sigma^2/2$. Approximating it crudely as a uniform sphere of radius R_e , we have $\mathcal{PE} = -3G\mathcal{M}^2/(5R_e)$ from Problem 3.12. Use Equation 3.44, the virial theorem, to show that the mass $\mathcal{M} \approx 5\sigma^2 R_e/G$. If all elliptical galaxies could be described by Equation 6.1 with the same value of n , explain why we would then have $\mathcal{M} \propto \sigma^2 R_e$ and the luminosity $L \propto I_e R_e^2$, so that $\mathcal{M}/L \propto \sigma^2/(I_e R_e)$.

(a) Show that, if all ellipticals had the same ratio \mathcal{M}/L and surface brightness $I(R_e)$, they would follow the Faber–Jackson relation.

(b) Show that Equation 6.19 implies that $I_e \propto \sigma^{1.5} R_e^{-1.25}$, and hence that $\mathcal{M}/L \propto \sigma^{0.5} R_e^{0.25}$ or $\mathcal{M}^{0.25}$: the mass-to-light ratio is larger in big galaxies.

unresolved (point) source

Surface brightness & flux:
unresolved sources



If source smaller than beam,
detect total flux of source

$$f_{\text{det}} = \int_{\Omega_{\text{beam}}} I \, d\Omega = \bar{I} \Omega_{\text{source}} \quad \text{if } \Omega_s \ll \Omega_b$$

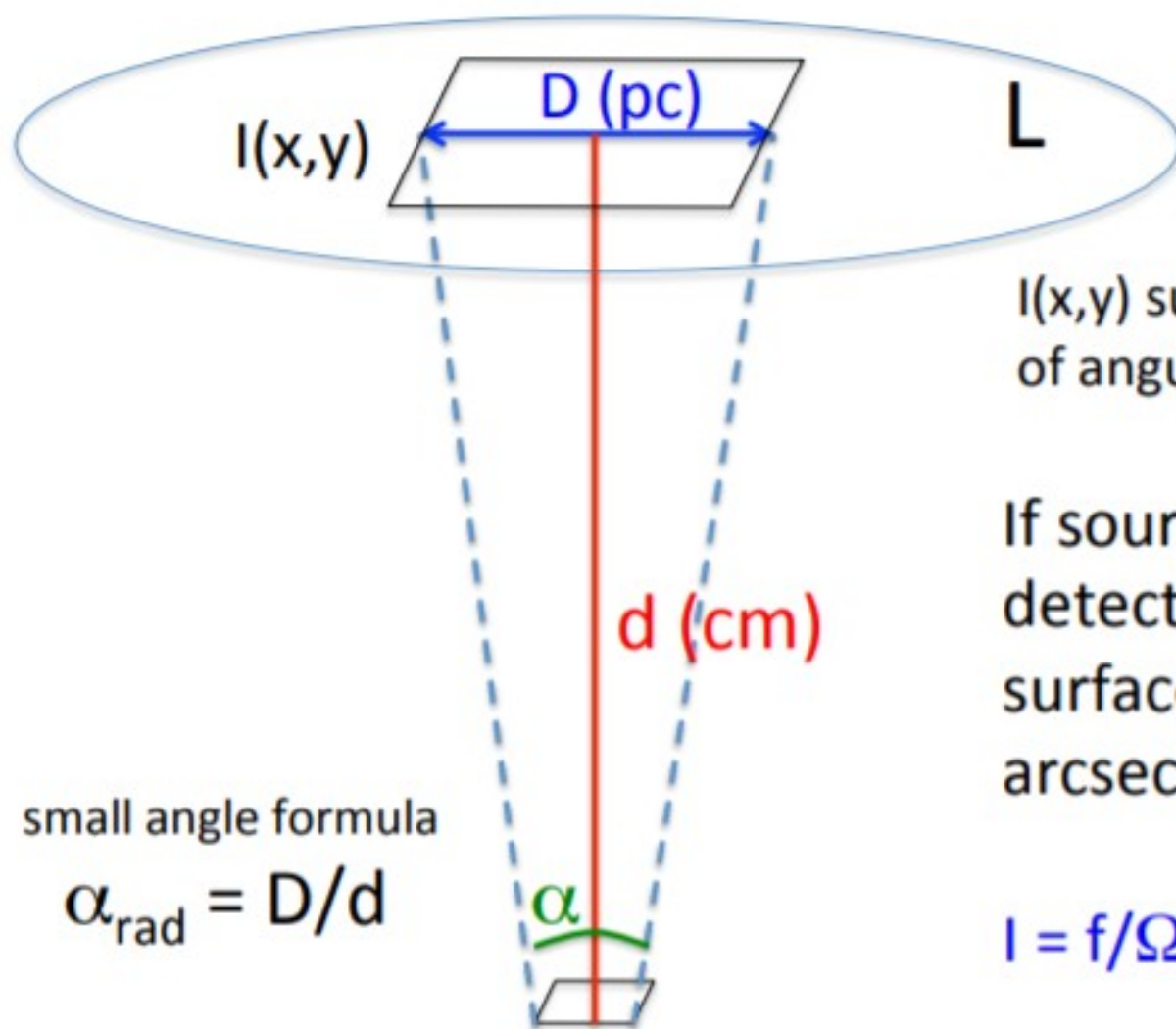
since $I(r) = 0$ for $r > r_s$

$$= f_{\text{tot}} = \frac{L}{4\pi d^2}$$

detects **all** of source flux

extended source

Surface brightness & flux:
resolved sources



$I(x,y)$ surface brightness I as a function of
of angular coordinates x,y

If source is **resolved**, a detector
detects the **flux per solid angle** =
surface brightness in $\text{erg s}^{-1} \text{cm}^{-2}$
 arcsec^{-2} (or sr^{-1})

small angle formula

$$\alpha_{\text{rad}} = D/d$$

$$I = f/\Omega = f/\alpha^2$$

solid angle
of square patch

$$\Omega = \alpha^2$$

detector
pixel

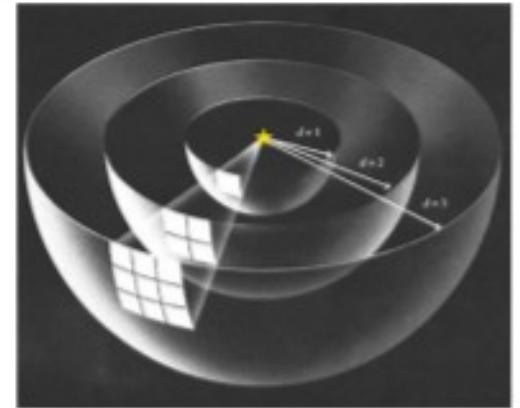
detects only part of source flux

Surface brightness is independent of distance!

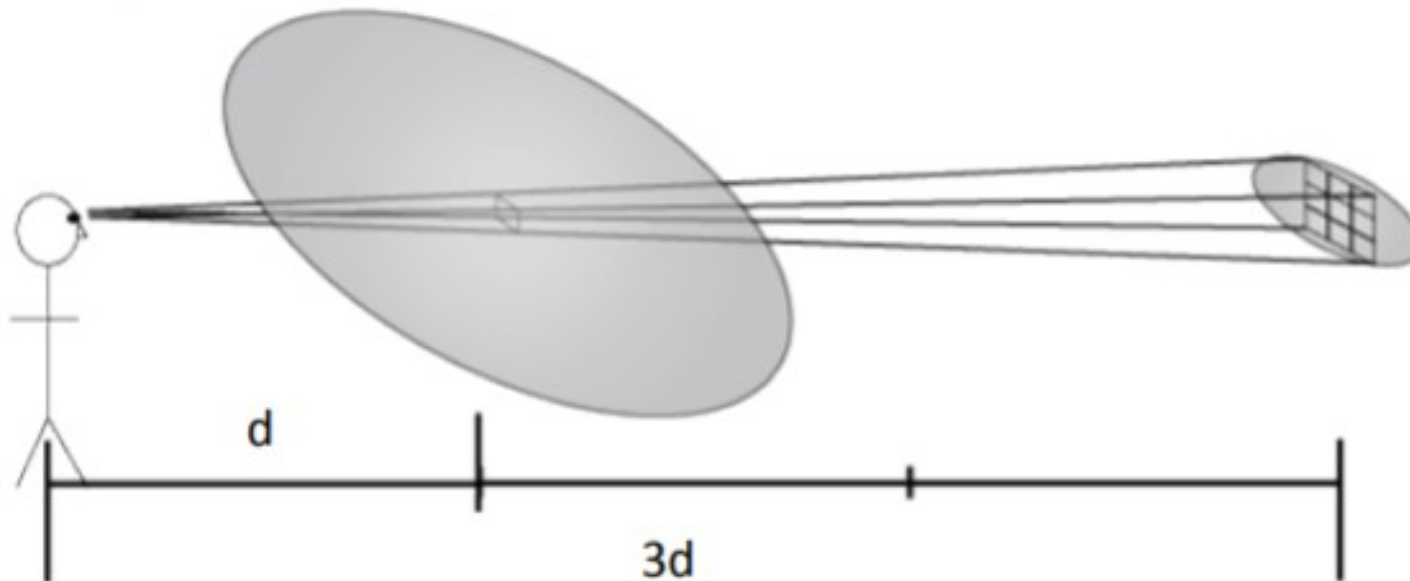
surface brightness = brightness or flux per solid angle

- Less light from each square meter of more distant source

(Inverse square law – B decreases by $1/d^2$)



- But more square meters (surface area) of source within same solid angle of observer for more distant source (surface area increases by d^2)



Surface brightness is distance independent

- If source is **unresolved**, a detector detects the **flux** in $\text{erg s}^{-1} \text{cm}^{-2}$
- If source is **resolved**, a detector detects the **flux per solid angle** = surface brightness in $\text{erg s}^{-1} \text{cm}^{-2} \text{arcsec}^{-2}$ (or sr^{-1})

$$I = f/\Omega = f/\alpha^2$$

Recall: angular size of source $\alpha = D/d$

angular area of source (square patch) $\Omega = \alpha^2 = (D/d)^2$

$f = L/4\pi d^2$ d = distance

$$I = f/\Omega = (L/4\pi d^2) / (D/d)^2 = L/4\pi D^2 \text{ where } D = \text{size of patch on source}$$

So units of I are $L_{\text{sun}} \text{pc}^{-2}$ or $\text{erg s}^{-1} \text{cm}^{-2} \text{arcsec}^{-2}$

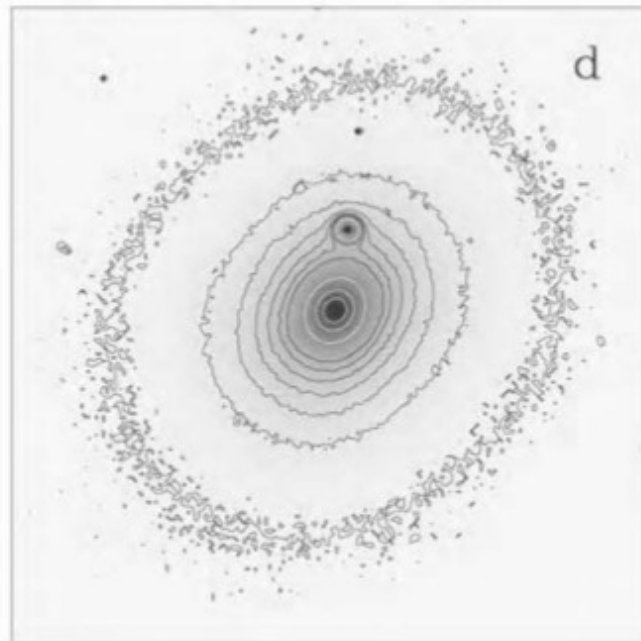
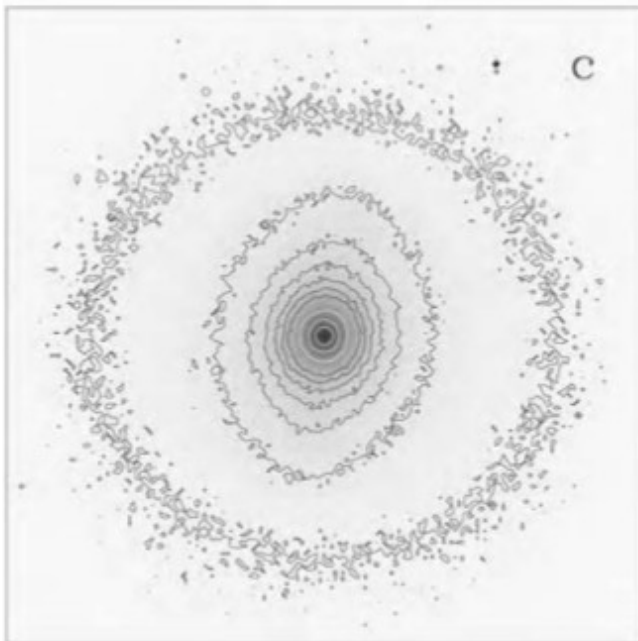
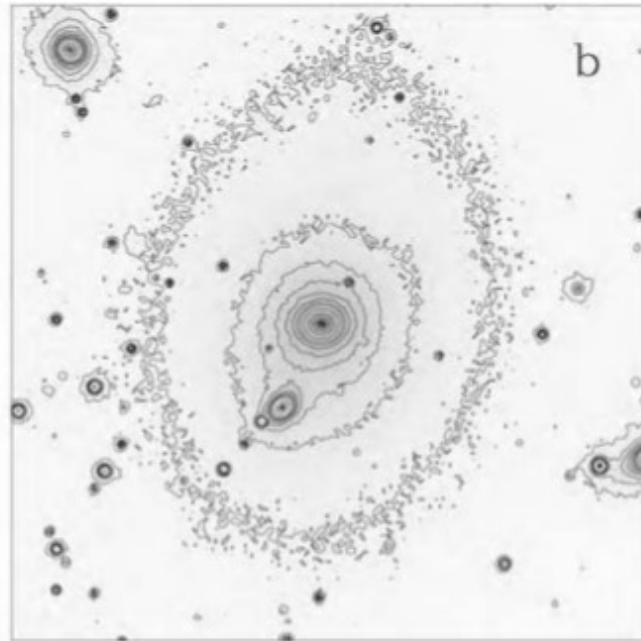
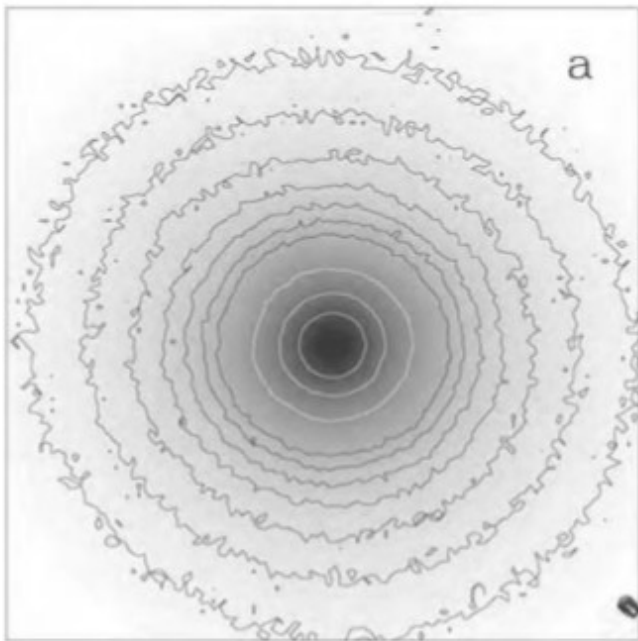
- Area on source (D^2 in pc^2) depends on distance (d in cm) and angular area ($\Omega = \alpha^2$ in arcsec^2); so that's why units of $\text{cm}^{-2} \text{arcsec}^{-2}$ are equivalent to pc^{-2} in SB
- Luminosity and area of patch in source both increase as d^2 so ratio doesn't depend on d !

Surface brightness in magnitudes arcsec⁻²

$$\mu = -2.5 \log I + C$$

SB in mag arcsec⁻² SB in erg s⁻¹ cm⁻² arcsec⁻²

- magnitudes arcsec⁻² are strange units since magnitudes are not linear: if a point in a galaxy has a SB of 21 magnitudes arcsec⁻² this means an area of 1 square arcsecond around this point emits as much light as a star of apparent magnitude 21.

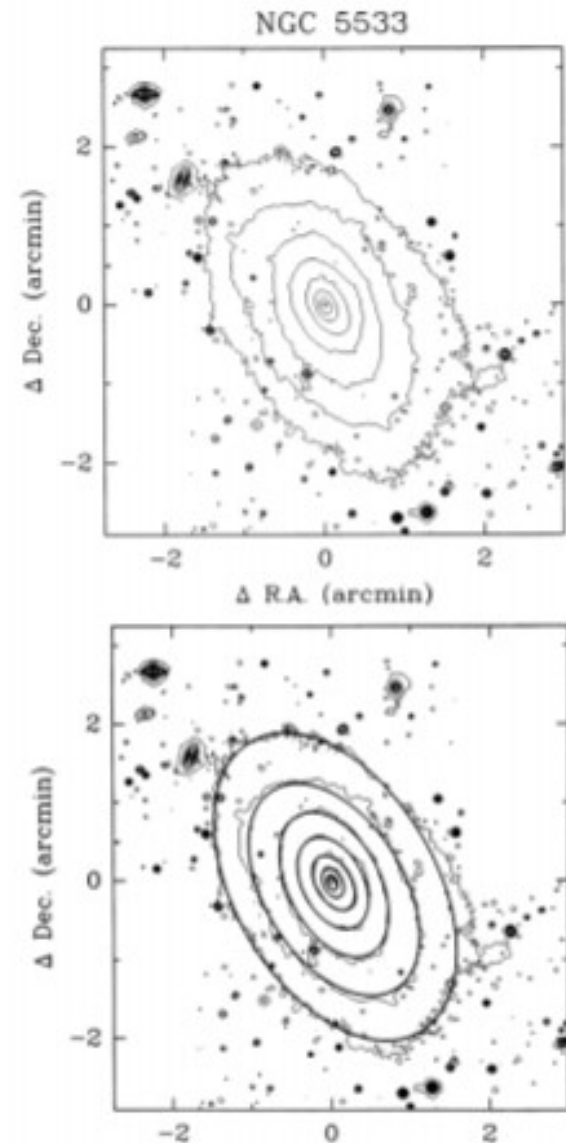


Color optical image of spiral galaxy



Separate images taken in 3 bands: g, r, i
3 images combined to make color image

Isophotes – contours of equal surface brightness



Fit ellipses to isophotes

Problem 6.13 The redshift of NGC 5266 is $cz \approx 3000 \text{ km s}^{-1}$; if $H_0 = 75 \text{ km s}^{-1} \text{ Mpc}^{-1}$, show that its distance $d \approx 40 \text{ Mpc}$. Use Equation 3.20 to show that the mass $\mathcal{M}(<4') \approx 7 \times 10^{11} \mathcal{M}_{\odot}$. The total apparent magnitude $B_{\text{T}}^0 = 12.02$; show that $L_B \approx 4 \times 10^{10} L_{\odot}$ – this is a big galaxy – so that the mass-to-light ratio $\mathcal{M}/L_B \approx 18$.

